

A combined *HST* and *XMM-Newton* campaign for the magnetic O9.5 V star HD 54879

Towards constraining the weak-wind problem of massive stars

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ABSTRACT

Context. HD 54879 (O9.7 V) is one of a dozen O-stars for which an organized atmospheric magnetic field has been detected. Despite their importance, little is known about the winds and evolution of magnetized massive stars.

Aims. To gain insights into the interplay between atmospheres, winds, and magnetic fields of massive stars, we acquired UV and X-ray data of HD 54879 using the *Hubble Space Telescope* and the *XMM-Newton* satellite. In addition, 35 optical amateur spectra were secured to study the variability of HD 54879.

Methods. A multiwavelength (X-ray to optical) spectral analysis is performed using the Potsdam Wolf-Rayet (PoWR) model atmosphere code and the *xspec* software.

Results. The photospheric parameters ($T_* = 30.5$ kK, $\log g = 4.0$ [cm s⁻²], $\log L = 4.45$ [L_\odot]) are typical for an O9.7 V star. The microturbulent, macroturbulent, and projected rotational velocities are lower than previously suggested ($\xi_{\text{ph}}, v_{\text{mac}}, v \sin i \leq 4$ km s⁻¹). An initial mass of $16 M_\odot$ and an age of 5 Myr are inferred from evolutionary tracks. We derive a mean X-ray emitting temperature of $\log T_X = 6.7$ [K] and an X-ray luminosity of $L_X = 1 \cdot 10^{32}$ erg s⁻¹. Short- and long-scale variability is seen in the H α line, but only a very long period of $P \approx 5$ yr could be estimated. Assessing the circumstellar density of HD 54879 using UV spectra, we can roughly estimate the mass-loss rate HD 54879 would have in the absence of a magnetic field as $\log \dot{M}_{B=0} \approx -9.0$ [$M_\odot \text{ yr}^{-1}$]. The magnetic field traps the stellar wind up to the Alfvén radius $r_A \gtrsim 12 R_*$, implying that its true mass-loss rate is $\log \dot{M} \lesssim -10.2$ [$M_\odot \text{ yr}^{-1}$]. Hence, density enhancements around magnetic stars can be exploited to estimate mass-loss rates of non-magnetic stars of similar spectral types, essential for resolving the weak wind problem.

Conclusions. Our study confirms that strongly magnetized stars lose little or no mass, and supplies important constraints on the weak-wind problem of massive main sequence stars.

Key words. Stars: Massive – Stars: magnetic field – Stars: mass-loss – Stars: individual: HD 54879

1. Introduction

Massive stars ($M_{\text{ini}} \gtrsim 8 M_\odot$) can outshine a million suns and radiate at energies that greatly exceed the ionization threshold of H, He I, and He II atoms. Through their powerful stellar winds and their final explosion as core-collapse supernova, they provide large amounts of kinetic energy and metals to their environment. Yet despite the importance of massive stars, our understanding of their evolution is prone to much debate. Along with binarity (e.g., Eldridge & Stanway 2016; Shenar et al. 2016, 2017), rotation (e.g., Georgy et al. 2012; Shenar et al. 2014; Shara et al. 2017), and metallicity effects (e.g., Crowther & Hadfield 2006; Hainich et al. 2015), another important question concerns the incidence of globally organized magnetic fields in massive stars and their impact on the stellar properties and evolution (see reviews by Donati & Landstreet 2009; Langer 2012).

Roughly 5 – 7% of the massive stars are estimated to possess global atmospheric magnetic fields (Wade et al. 2014; Schöller et al. 2017; Grunhut et al. 2017), the majority of which are B-type.

The first detection of an organized magnetic field in an O-type star was reported by Donati et al. (2002) for θ^1 Orionis C. Since then, about ten O-stars were added to the sample, thanks to the Magnetism in Massive Stars (MiMeS: Grunhut et al. 2009; Petit et al. 2011; Alecian et al. 2014; Wade et al. 2016), B fields in OB stars (BOB: Morel et al. 2014; Fossati et al. 2015; Hubrig et al. 2015a), and magnetic field origin (MAGORI: Hubrig et al. 2011) collaborations. It is commonly believed that massive stars may become magnetic if they originate in significantly magnetized molecular clouds, and hence their magnetic fields are fossil. Alternatively, global magnetic fields may form as a result of merger events or dynamos produced during the pre-main sequence phase (Moss 2003; Ferrario et al. 2009; Donati & Landstreet 2009).

Atmospheric magnetic fields can impact surface rotation rates via magnetic braking (Weber & Davis 1967; Ud-Doula et al. 2008), introduce chemical abundance inhomogeneities and peculiarities (Hunger & Groote 1999), and confine the stellar wind in a so-called magnetosphere (e.g., Friend & MacGregor 1984; ud-Doula & Owocki 2002; Townsend et al. 2005). As a result of

the latter, magnetic fields can significantly reduce the mass-loss from the star, favoring the creation of massive compact objects upon core collapse (e.g., Petit et al. 2017). As the ejected matter streams along the field lines towards the magnetic equator, powerful collisions occur that produce copious X-ray emission (Babel & Montmerle 1997). Considering how little is known about the incidence, evolution, and impact of magnetic fields, studying massive magnetized stars is essential.

The subject of our study, HD 54879, was classified as O9.7 V by Sota et al. (2011). Castro et al. (2015, C2015 hereafter) measured a longitudinal magnetic field reaching a maximum of $|B_z| \approx 600$ G for HD 54879, from which they estimated a dipole field of $B_d \geq 2$ kG. The star is believed to reside in the CMa OB1 association, for which an age of ≈ 3 Myr was estimated (Clariá 1974). There are several distance estimates for the CMa OB1 cluster (e.g., Clariá 1974: 1.15 ± 0.14 kpc, Humphreys 1978: 1.32 kpc; Kaltcheva & Hilditch 2000: 0.99 ± 0.05 kpc). The first data release by the *Gaia* satellite gives $d = 0.86^{+0.36}_{-0.20}$ kpc for HD 54879 (Gaia Collaboration et al. 2016). Following Gregorio-Hetem (2008), we adopt $d = 1.0 \pm 0.2$ kpc in this study.

C2015 showed that a very large mass-loss rate ($\log \dot{M} \approx -5.0$ [$M_\odot \text{ yr}^{-1}$]) is necessary to reproduce the $H\alpha$ emission of HD 54879. This value is orders of magnitude larger than what is expected for late O-type dwarf stars. C2015 therefore suggested that the $H\alpha$ emission originates in the magnetosphere. In fact, late massive main sequence stars are generally known to exhibit mass-loss rates which are significantly lower than predicted by theory (Vink et al. 2000), often referred to as the *weak wind problem* (e.g., Marcolino et al. 2009).

Our study benefits from new UV and X-ray spectra acquired by us simultaneously with the *Hubble Space Telescope* (*HST*), and the *XMM-Newton* satellite (see Sect. 2 for details). Complemented by optical HARPS spectra, the data at hand allow for a multiwavelength spectral analysis of the stellar photosphere and magnetosphere of HD 54879 (Sect. 3), performed here using the Potsdam Wolf-Rayet (PoWR) code. The evolutionary channel of HD 54879 is discussed in Sect. 4. X-ray data are analyzed using the *xSPEC* software (Sect. 5). The combination of X-ray, UV, and optical data is essential, as these spectral ranges probe different regions of the magnetosphere and stellar wind (see Fig. 1 and Sect. 3.3). Additionally, we collected 35 amateur spectra covering the $H\alpha$ line to study the spectral variability of the star (Sect. 6). A follow up study (Järvinen et al. in prep.) will focus on the structure and variability of the magnetic field of HD 54879.

2. Observational data

We acquired three UV spectra of HD 54879 on April 30, 2016 (PI: Hamann, proposal ID: 14480) using the *HST* high resolution ($R = 45\,800$) Space Telescope Imaging Spectrograph (STIS). The spectra cover the spectral range 1123 – 1710 Å. As no notable variability in the spectral lines could be identified, we co-added the three exposures to obtain a single spectrum with $\text{SNR} \approx 100$. While the exposures are flux calibrated, they have a significant offset relative to one another of the order of a factor of two, likely associated with light losses originating in thermal “breathing” that changes the focus over orbital time scales (Proffitt et al. 2017). The co-added spectrum was therefore recalibrated to match photometry from the TD-1 satellite (Thompson et al. 1978).

XMM-Newton observed HD 54879 on May 01, 2016 (PI: Hamann, ID: 0780180101) with an exposure time of ≈ 40 ks. All three European Photon Imaging Cameras (EPICs: MOS1, MOS2, and PN) were operated in the standard, full-frame mode with the medium UV filter. The observations were affected by periods of

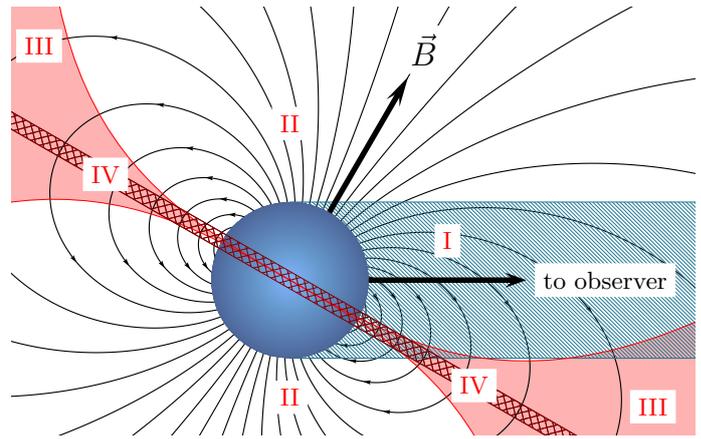


Fig. 1: A schematic sketch of a star with a global dipole magnetic field, illustrating the formation regions of different features in the spectrum of HD 54879. Region I: blue-shifted UV resonance line absorption (e.g., C iv, Si iv, N v); Region II: UV resonance line emission; Region III: shocked, X-ray emitting region; Region IV: recombination line emission at the magnetic equator (e.g., $H\alpha$, $H\beta$).

high flaring background, likely caused by soft protons populating Earth’s magnetosphere¹. After excluding these periods, the useful exposure time was reduced to 29.8 ks. The data were reduced using the most recent calibration. The spectra and light-curves were extracted using standard procedures from a region with a diameter of about $15''$. The background area was chosen to be nearby the star and free of X-ray sources.

Three spectropolarimetric observations of HD 54879 were obtained with the HARPS polarimeter (HARPSpol, Snik et al. 2008) attached to ESO’s 3.6 m telescope (La Silla, Chile) within ESO Large programme ID 191.D-0255 (PI Morel) on April 22, 2014 as well as on March 11 and 14, 2015. In this study, we make use only of the intensity spectra, which are of high-resolution ($R = 115\,000$), high signal-to-noise ratio ($\text{SNR} \approx 300$), and cover the wavelength range 3780 – 6912 Å with a gap at 5259 – 5337 Å.

To study the spectral variability, we employed 35 spectra from the Shenton Park Observatory (SPO) taken by Paul Luckas using a 0.35m Ritchey-Chrétien telescope equipped with a Shelyak Lhires III spectrograph operating at a resolution of $R \approx 16\,000$ and producing spectra with $\text{SNR} \approx 100$. The spectra cover the spectral range 6500 – 6610 Å ($H\alpha$). Spectral images were bias, dark, and flat field corrected in the normal manner, and calibrated using Ne/Ar arc lamp spectra taken nightly and adjacent to science imaging. We also use seven spectra taken with the FEROS and FIES spectrographs from the IACOB and OWN projects (Barbá et al. 2010; Simón-Díaz et al. 2011a,b; Simón-Díaz & Herrero 2014) between the years 2009 and 2013 (see latter references for details). Together with the HARPS data, this makes a total of 45 spectra used to study the spectral variability. The spectra were cleaned from tellurics using the ESO tool *Molecfit* (Smette et al. 2015; Kausch et al. 2015) and rectified by eye. A log containing the epochs of observation is given in Table A.1 in Appendix A.

For the spectral analysis, we also use UV photometry taken by the TD-1 satellite (Thompson et al. 1978), *UBV* photometry from Myers et al. (2003), *R* and *JHK_s* 2MASS photometry from Zacharias et al. (2005), and *I*-band photometry from Monet et al. (2003).

¹ see *XMM-Newton Users Handbook*

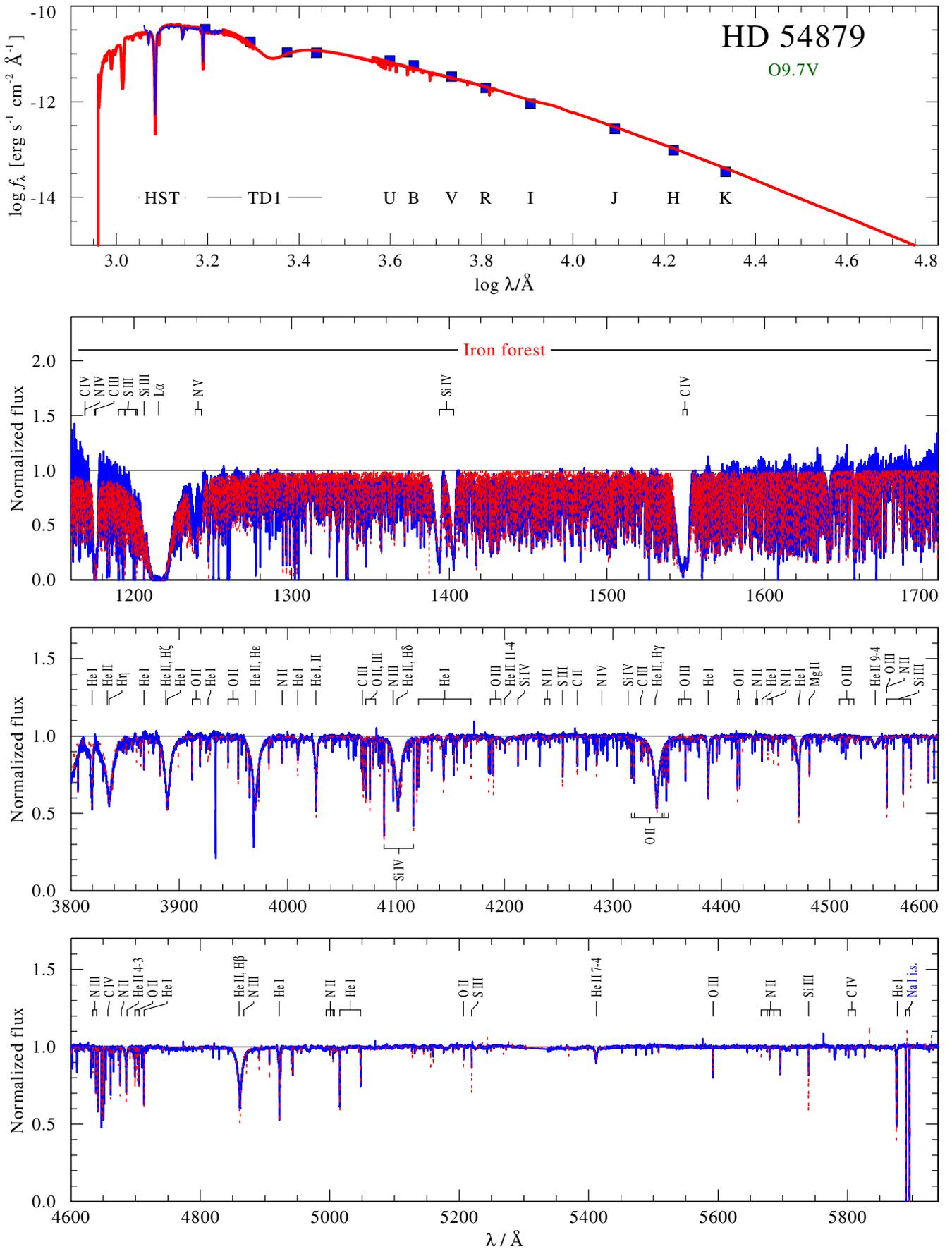


Fig. 2: *Upper panel*: Comparison between observed photometry (blue squares) and the flux-calibrated *HST* spectrum (blue line) with the SED of our best-fitting model (red line). *Lower panels*: Comparison between normalized *HST* and HARPS spectra (blue solid line) and the best-fitting model (red dotted line). H α is shown separately (see Fig. 4)

3. Spectral analysis

With the goal of inferring the fundamental stellar and wind parameters, we now perform a consistent, multiwavelength analysis of the optical and UV spectra at hand. HD 54879 was spectroscopically analyzed before (C2015), but the analysis did not include UV data. The spectral modelling is performed with the PoWR model atmosphere code². The PoWR code solves the radiative transfer and statistical equations in an expanding, spherically-symmetric atmosphere, relaxing the assumption of local thermodynamic equilibrium (i.e., non-LTE). Indeed, the spherical symmetry in HD 54879 is broken by the presence of the strong magnetic dipole. A complete non-LTE solution of a magnetized stellar atmosphere in multiple dimensions is currently unfeasible. Nevertheless, our model provides a good approximation for the stellar photosphere, and can deliver significant insights on the stellar wind and mass-loss, as we show in the following sections. For more details on the PoWR code, we refer to Gräfener et al. (2002) and Hamann & Gräfener (2003).

By fitting synthetic spectra to the observations, we derive the effective temperature T_* , the surface gravity g_* , and the stellar luminosity L . The effective temperature T_* refers to the stellar radius R_* , so that $L = 4\pi\sigma R_*^2 T_*^4$. The stellar radius R_* is defined at the model's inner boundary, fixed at a mean Rosseland optical depth of $\tau_{\text{Ross}} = 20$. The velocity field consists of two regimes. In the subsonic regime, hydrostatic equilibrium is approached (Sander et al. 2015). In the supersonic regime, the velocity follows the β -law with the value $\beta = 0.8$, typical for O-stars (Castor et al. 1975; Kudritzki et al. 1989; Puls et al. 1996). The co-moving frame radiative transfer is calculated adopting Gaussians for the absorption/emission coefficients with a constant Doppler width of $v_{\text{Dop}} = 20 \text{ km s}^{-1}$. During the calculation of the emergent spectrum, v_{Dop} is calculated via $v_{\text{Dop}} = (v_{\text{th}}^2 + \xi^2)^{1/2}$, where v_{th} is the thermal velocity, and the microturbulent velocity $\xi(r)$ is assumed to grow from the photospheric value ξ_{ph} to the peak value ξ_{max} at a prespecified radius $R_{\xi_{\text{max}}}$ (see detailed in Shenar et al. 2015). A depth-dependent wind clumping is assumed here, reaching a maximum of $D = 10$. The synthetic profiles are convolved with Gaussians corresponding to the respective instrumental profiles. A comparison between the best-fitting model and the observations is shown in Fig. 2. The derived stellar parameters are given in Table 1.

3.1. The stellar photosphere

Figure 3 shows several photospheric lines and the $\text{H}\alpha$ line as observed in the three available HARPS spectra (see Table A.1). As the figure illustrates, the $\text{H}\alpha$ line varies notably in shape, width, and equivalent width (EW). The intriguing three-peak profile of $\text{H}\alpha$ is persistent in all observations. In contrast, no variability is detected in the photospheric features in the HARPS spectra. In fact, within the limitation set by the spectral resolution and SNR, no notable variability of the photospheric features is detected in the FEROS, FIES, or HARPS spectra taken between the years 2009 and 2015. He II lines also show no evidence for variability, unlike in other massive magnetic stars (e.g., Grunhut et al. 2012; Hubrig et al. 2015b). Assuming that the star changes its orientation over the span of six years because of its rotation, this seems to suggest that the spherical symmetry of the photospheric layers is virtually unbroken by the presence of the magnetic field. Until evidence for photospheric variability is observed in HD

54879, a spherical symmetric model appears to be adequate for the analysis of its photosphere.

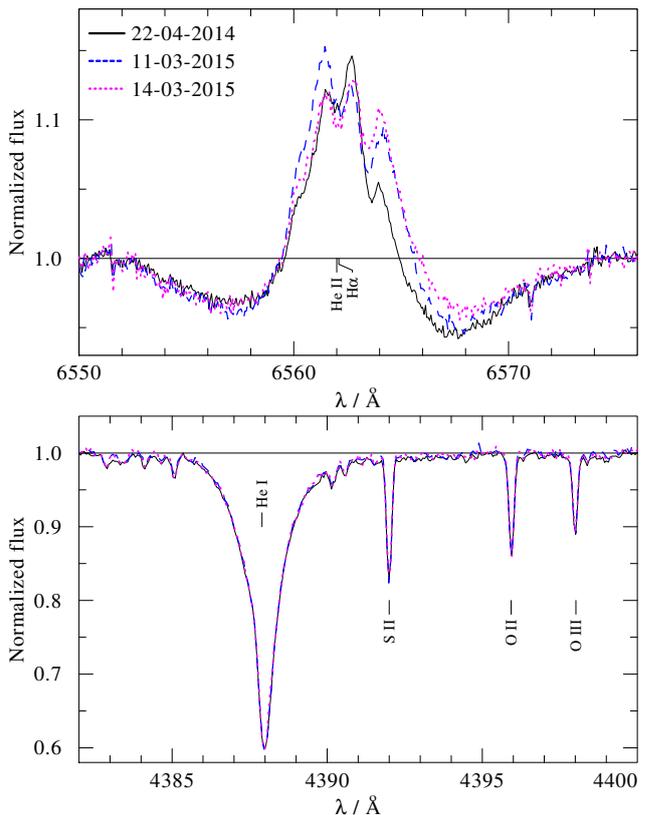


Fig. 3: The three HARPS observations focusing on $\text{H}\alpha$ (upper panel) and a few photospheric lines (lower panel).

The surface gravity $\log g_*$ is derived primarily from the pressure broadened wings of the $\text{H}\gamma$ and $\text{H}\delta$ lines as well as He II lines in the optical spectrum (see Fig. 4). $\text{H}\alpha$, and to a lesser extent $\text{H}\beta$, are strongly contaminated by emission originating at the magnetic equator (Sundqvist et al. 2012; ud-Doula et al. 2013). Since this is an inherently non spherically symmetric structure (see Fig. 1 and Sect. 3.3), we do not attempt to include the disk emission in our model, and ignore $\text{H}\alpha$ and $\text{H}\beta$ for the estimation of $\log g_*$. Our estimate for $\log g_*$ is found to be consistent with the value reported by C2015.

The temperature is determined from the ratios between lines from different ionization stages of the same element, e.g., He I/II (see Fig. 5), N II/III/IV, and C II/III/IV. The temperature we obtain ($30.5 \pm 0.5 \text{ kK}$) is significantly lower than that derived by C2015 ($33 \pm 1 \text{ kK}$). We note, however, that the temperature derived here is more consistent with our target's spectral type, O9.7 V, considering calibrations by, e.g., Martins et al. (2005a). Moreover, a comparison of the observed photospheric spectra with TLUSTY model atmospheres (Hubeny & Lanz 1995; Lanz & Hubeny 2003) implies a similar temperature to that derived here.

By fitting the star's spectral energy distribution (SED), we derive the luminosity L and reddening E_{B-V} (see upper panel of Fig. 2). Different extinction laws were tested, and the SED is best reproduced using laws published by Seaton (1979) and Nandy et al. (1975). The luminosity derived here agrees well with C2015 when scaling $L \propto d^2$ to the distance used by C2015 ($d = 1.32 \text{ kpc}$), which is likely overestimated (see discussion on the distance in Sect. 1).

² PoWR models of Wolf-Rayet and OB-type stars can be downloaded at www.astro.physik.uni-potsdam.de/PoWR

Table 1: Derived physical parameters for HD 54879

T_* [K]	$\log L$ [L_\odot]	$\log g$ [cm s^{-1}]	R_* [R_\odot]	M_* [M_\odot]	$v \sin i$ [km s^{-1}]	ξ_{ph} [km s^{-1}]	v_{mac} [km s^{-1}]	$\log \dot{M}_{B=0}^b$ [$M_\odot \text{ yr}^{-1}$]	$\log \dot{M}$ [$M_\odot \text{ yr}^{-1}$]	E_{B-V} [mag]
30.5 ± 0.5	4.45 ± 0.2^a	4.0 ± 0.1	6.1 ± 1.5^a	14 ± 7^a	< 4	< 4	< 4	≈ -9.0	$\lesssim -10.2$	0.35 ± 0.01

Notes. Stellar parameters derived from our spectral analysis.

^(a) The errors are dominated by the assumed uncertainty in the distance ($\approx 20\%$, see Sect. 1).

^(b) The mass-loss rate that the star would have in the absence of a magnetic field should not be confused with the true mass-loss rate.

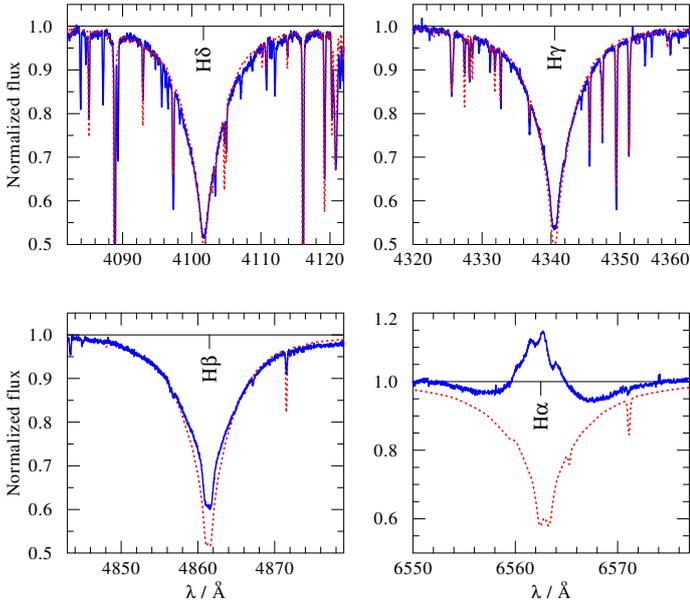


Fig. 4: Our best fitting model with parameters as in Table 1 (red dotted line) compared to the first four Balmer lines, as observed in the HARPS spectrum (blue solid line). The disk emission of $H\alpha$ and $H\beta$ is not included in the model.

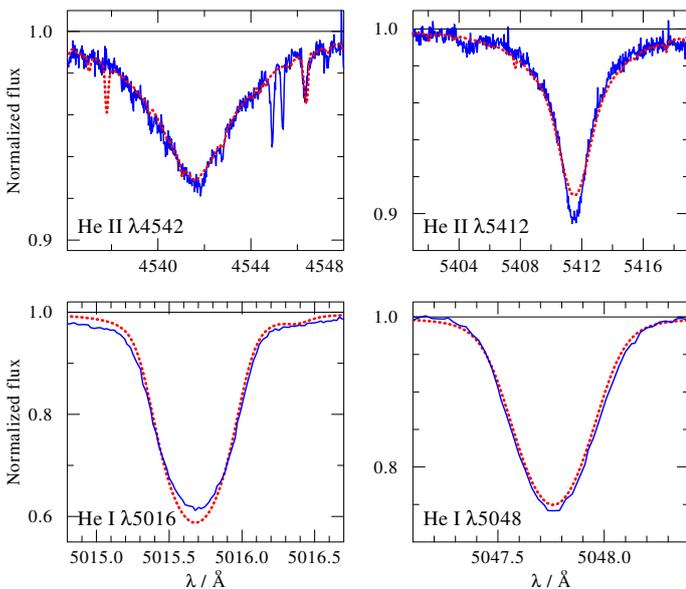


Fig. 5: Same as Fig. 4, but for several He I, II lines.

Finally, having constrained the fundamental parameters of the star, the abundances are inferred from the overall strength of spec-

tral lines belonging to the corresponding element. Abundances that could not be derived due to the absence of corresponding spectral lines are assumed to have solar values (Asplund et al. 2009). The derived/adopted values are given in Table 2. We can exclude a significant overabundance of He compared to the solar value. Our results agree well with those by C2015, within errors. It is interesting to note that both C and N are found to be significantly subsolar, standing in contrast to reports of nitrogen enhancement in magnetic B-type stars (e.g., Morel et al. 2008).

3.2. Rotation and turbulence

Any periodicities observed in our target are likely to correlate with its rotational period P_{rot} . It is therefore important to constrain the projected rotational velocity $v \sin i$ from our high resolution spectra. C2015 used the Fourier tool *IACOB-BROAD* (Simón-Díaz & Herrero 2014) to derive a projected rotational velocity of $v \sin i = 7 \pm 2 \text{ km s}^{-1}$. However, for rotational velocities that are comparable to the microturbulent velocity, the tool strongly underestimates the involved errors, making it inapplicable for HD 54879 (see discussion by Simón-Díaz & Herrero 2014, Sect. 3.4). Constraining $v \sin i$ is difficult, because its associated spectral line broadening “competes” with the microturbulent velocity ξ_{ph} , the macroturbulent velocity v_{mac} , and the thermal velocity $v_{\text{th}} = 2k_{\text{B}}T/m$. To overcome this, we consider first spectral lines belonging to CNO ions. For these elements, $v_{\text{th}} \approx 6 - 7 \text{ km s}^{-1}$, which alone agrees with the observed line widths. Therefore, the unknowns ξ_{ph} , $v \sin i$, and v_{mac} must be smaller than v_{th} .

Figure 6 displays selected CNO spectral lines compared to our best fitting model convolved with rotation profiles of different $v \sin i$ values (see legend and caption) calculated with $\xi_{\text{ph}} = 0 \text{ km s}^{-1}$. It is evident that, for the majority of spectral lines, no broadening mechanism beyond those which are included in the calculation are required. Figure 6 shows that we can only hope to derive an upper limit for the projected rotation velocity. Based on our study, we can constrain this limit to $v \sin i \leq 4 \text{ km s}^{-1}$.

The microturbulence could only be estimated approximately from the UV iron lines, which are subject to low thermal broadening ($v_{\text{th}} \approx 3 \text{ km s}^{-1}$). We estimate $\xi_{\text{ph}} \approx 3 \text{ km s}^{-1}$, but can only reliably constrain an upper limit of $\xi_{\text{ph}} \leq 4 \text{ km s}^{-1}$. Similarly, we can conclude $v_{\text{mac}} \leq 4 \text{ km s}^{-1}$, as opposed to $v_{\text{mac}} = 8 \text{ km s}^{-1}$ reported by C2015.

The small value of v_{mac} derived here stands in contrast to the significant macroturbulent velocities reported for other magnetic O-type stars (Sundqvist et al. 2013). A very low macroturbulent velocity ($< 3 \text{ km s}^{-1}$) was reported by the latter authors only for the strongly magnetized star NGC 1624-2. However, Sundqvist et al. (2013) showed that this can be understood as a consequence of its extraordinary magnetic field ($B_{\text{d}} = 20 \text{ kG}$), which is strong enough to stabilize the atmosphere at deep sub-photospheric layers, where the iron opacity peak is reached and macroturbulence

Table 2: Derived chemical abundances (in mass fractions) for HD 54879

	X_{H}	X_{He}	$X_{\text{C}}/10^{-3}$	$X_{\text{N}}/10^{-4}$	$X_{\text{O}}/10^{-3}$	$X_{\text{Ne}}/10^{-3}$	$X_{\text{Mg}}/10^{-4}$	$X_{\text{Al}}/10^{-5}$	$X_{\text{Si}}/10^{-4}$	$X_{\text{P}}/10^{-6}$	$X_{\text{S}}/10^{-4}$	$X_{\text{Fe}}/10^{-3}$
derived value	0.74	0.25	1 ± 0.3	3 ± 1	6 ± 2	1.26	7 ± 3	4 ± 2	7 ± 3	5.83	3 ± 1	1.2 ± 0.5
relative to solar	1	1	0.4	0.45	1	1	1	0.8	1	1	1	1

Notes. Values without errors were not derived and are assumed to be solar.

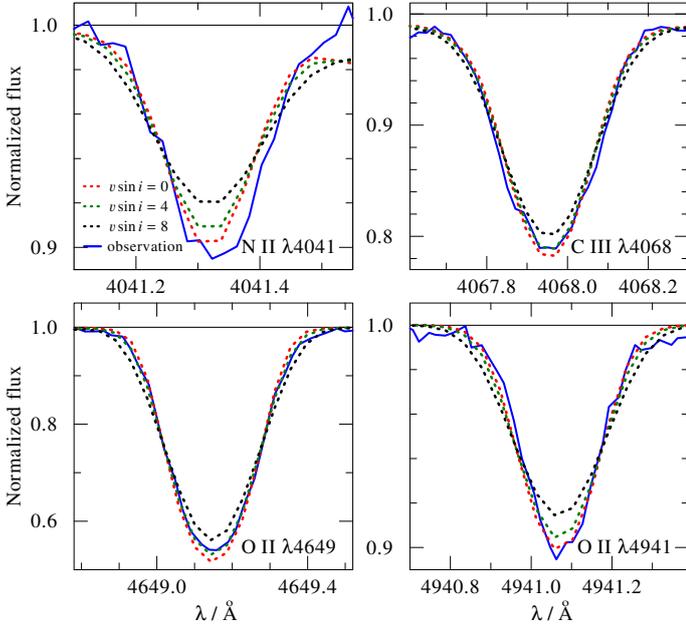


Fig. 6: Selected CNO spectral lines of our best fitting model convolved with rotation profiles for $v \sin i = 0, 4,$ and 8 km s^{-1} (red, green, and black dotted lines, respectively), compared compared to observations (blue solid line).

is thought to originate (e.g., Cantiello et al. 2009). The same argument does not hold for our target given its much weaker magnetic field. Rather, we suggest that the small value found here for HD 54879 is related to its low effective temperature (compared to the sample analyzed by Sundqvist et al. 2013). τ Sco is another magnetic star of a similar spectral type and magnetic field strength for which a low macroturbulent velocity was reported (Smith & Karp 1978). Recent studies by Simón-Díaz et al. (2017) imply that late-type massive stars exhibit systematically lower macroturbulent velocities compared to early-type stars, although the large scatter in v_{mac} values prevents us from concluding this unambiguously.

3.3. The stellar wind

Empirically derived mass-loss rates of low luminosity ($\log L_{\text{bol}}/L_{\odot} < 5.2$) OB-dwarfs are orders of magnitude lower than predicted by standard mass-loss recipes (Bouret et al. 2003; Martins et al. 2005b; Marcolino et al. 2009; Oskinova et al. 2011; Huenemoerder et al. 2012). This is often referred to in the literature as the *weak wind problem*.

There are various explanations as to why measured mass-loss rates of OB-dwarfs are so low. Some authors suggest that the wind-driving force in OB dwarfs is lower than predicted by current models, and that their winds are genuinely weak. For

example, Drew et al. (1994) proposed that the ionization of winds by X-rays reduces the total radiative acceleration. However, Oskinova et al. (2011) showed that ionization by X-rays does not significantly inhibit the wind driving power in magnetic B-dwarfs. An alternative idea was proposed by Lucy (2012), who suggested that, in late-type O dwarfs, the shock-heating of the ambient gas results in a single-component flow with a temperature of a few MK. This hot wind coasts to high velocities as a pure coronal wind. Hence, the bulk of the wind is not visible in the optical/UV, and mass-loss rates derived from these spectral ranges are not reliable. High-resolution X-ray spectroscopy of O-dwarfs seem to support this scenario (e.g., Huenemoerder et al. 2012, and references therein).

A comparison of *HST* UV data of the prototypical O9.7 V star ν Ori (HD 36512) with the *HST* data collected for our target, HD 54879, is shown in Fig. 7, where we focus on the two most prominent wind features, the resonance lines C IV $\lambda\lambda 1548, 1551$ and Si IV $\lambda\lambda 1394, 1403$. It is evident that ν Ori shows no, or very little, evidence for a stellar wind. Assuming ν Ori has the same fundamental parameters as HD 54879, and adopting a terminal velocity of $v_{\infty} = 1700 \text{ km s}^{-1}$ (Kudritzki & Puls 2000), a quick analysis of its UV spectrum using the PoWR code suggests an upper limit of $\log \dot{M} \lesssim -9.3 [M_{\odot} \text{ yr}^{-1}]$ for ν Ori. This is lower by almost two orders of magnitude compared to predictions by Vink et al. (2000), following the trend of other weak-wind stars.

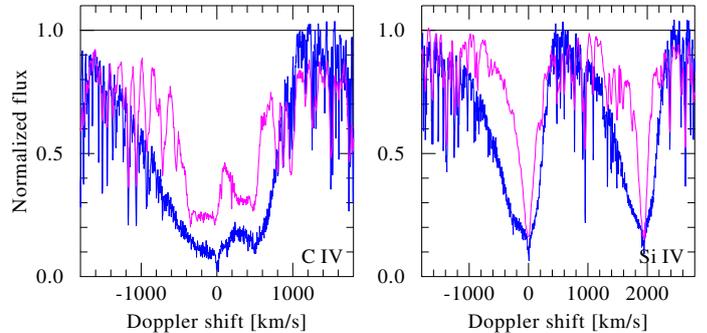


Fig. 7: Comparison of normalized *HST* observations of the prototypical O9.7 V star HD 36512 (ID: 13346, PI: Ayres, pink solid line) and of HD 54879 (blue solid line) in Doppler space. Shown are the resonance lines C IV $\lambda\lambda 1548, 1551$ (left panel) and Si IV $\lambda\lambda 1394, 1403$ (right panel).

In contrast to ν Ori, our target shows a clear asymmetry that is suggestive of absorption stemming from matter surrounding the star. At first glance, this seems to suggest that the magnetic field in fact spurs mass-loss from the star, contrary to expectation. However, the profiles of the resonance lines of HD 54879 have unusually shallow blueshifted edges, suggesting the presence of a large velocity dispersion, reaching a maximum speed of $\approx 1000 \text{ km s}^{-1}$. Similar line profiles were reported in other studies of magnetic stars (e.g., Marcolino et al. 2013; Nazé et al. 2015).

The magnetic field in HD 54879 dominates the behavior of the stellar wind for radii smaller than the Alfvén radius r_A , which is estimated to $r_A \gtrsim 12 R_*$ in this work (see below). To understand the line profiles of the Si IV and C IV resonance lines (Fig. 7), as well as the prominent $H\alpha$ emission (Fig. 4), we refer the reader to Fig. 1. The absorption of the UV resonance lines is mostly blueshifted, because it is formed in front of the stellar disk, where the material moves mostly towards the observer (region I). Some emission in the resonance lines (blueshifted and redshifted) is formed in the less dense region II, but it competes with line absorption originating in the stellar photosphere. Moreover, X-rays formed in region III (Owocki et al. 2016) ionize the Si IV and C IV ions and reduce their emissivity/opacity. Finally, as a recombination line, the emission in $H\alpha$ scales with ρ^2 , ρ being the density. Therefore, $H\alpha$ emission stems primarily from the magnetic equator (region IV), where ρ is orders of magnitudes larger. In case of a pronounced shock retreat, the formation of a disk-like structure may be inhibited (ud-Doula et al. 2014). However, $H\alpha$ is still expected to trace dense environments that strongly deviate from spherical symmetry, and is therefore not modelled in the framework of this study. Both the dynamic nature of the disk as well as possible co-rotation contribute to the width of the $H\alpha$ feature. The presence of $H\alpha$ emission or broadened, asymmetric lines in the UV is therefore not related to the stellar wind, but to the magnetosphere, and they do not imply a mass-loss from HD 54879 (see also Nazé et al. 2015; Erba et al. 2017).

To illustrate the inability of a standard, spherically-symmetric outflow to reproduce the observations, we plot in Fig. 8 several synthetic spectra calculated with the parameters given in Table 1, but with $v_\infty = 900 \text{ km s}^{-1}$ (roughly corresponding to the observed blue edge of the lines), a standard wind turbulent velocity ($0.1 v_\infty$), and different mass-loss rates. It is evident from the figure that no combination of wind parameters can reproduce the observed shallow profiles, which indicates that the stellar wind is heavily affected by the magnetic field.

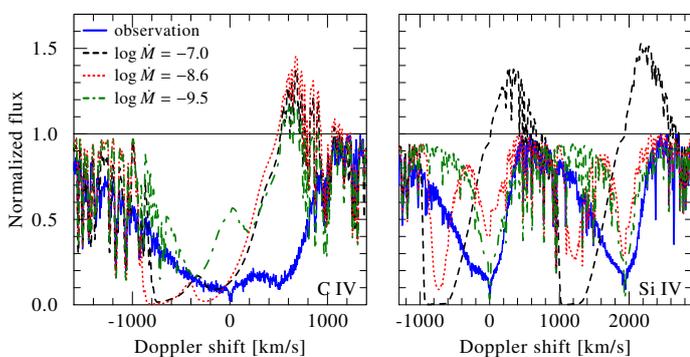


Fig. 8: Normalized *HST* observations of HD 54879 (blue solid line) compared to synthetic spectra with a terminal velocity of $v_\infty = 900 \text{ km s}^{-1}$ and different mass-loss rates (given in the legend in $[M_\odot \text{ yr}^{-1}]$).

Solving the full non-LTE radiative transfer in the presence of a magnetic field in 3D is currently not feasible and beyond the scope of this study. Instead, we simulate the motion of the matter along the field lines by a turbulent velocity which is strongly enhanced close above the stellar surface, at $r \approx 1.1 R_*$. We find the best fit for the combination of $v_{\infty, \text{sph}} = 300 \text{ km s}^{-1}$, $\xi_{\text{wind, sph}} = 500 \text{ km s}^{-1}$, and $\log \dot{M}_{\text{sph}} = -8.8 [M_\odot \text{ yr}^{-1}]$. X-rays ionization is approximately accounted for in a spherically-symmetric fashion (see Sect. 3.4). The results are shown in Fig. 9. Note that these parameters do not correspond to actual physical

parameters, since they are obtained via a spherically symmetric model which assumes a constant outflow. Rather, these parameters serve to reproduce the conditions in the formation region of the resonance lines. For example, the model predicts that the density at the main line forming region, at $r \approx 1.1 - 2 R_*$, is $\rho \approx 10^{-15} - 10^{-16} \text{ g cm}^{-3}$.

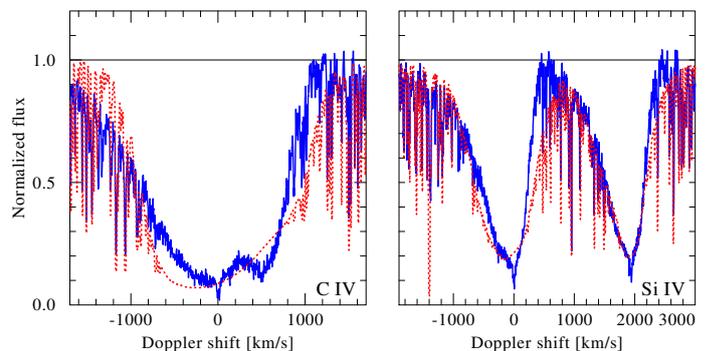


Fig. 9: Normalized *HST* observations of HD 54879 (blue solid line) compared to our best-fitting model, which accounts for a large microturbulence and superionization by X-rays.

By comparing the derived densities to predictions by an analytical model derived for stars with a constant outflow and a global dipole magnetic field (Owocki et al. 2016), we can give a rough estimate to the mass-loss rate HD 54879 would have in the absence of a magnetic field, $\dot{M}_{B=0}$. This model assumes three components: (a) an upflow along the dipole lines towards the loop apex, (b) a shocked, X-ray emitting plasma, and (c) a cooled downflow from the dipole apices back towards the star. Assuming simplistically a β -type law for the expansion speed, the model predicts the densities ρ_u , ρ_s , and ρ_d (Eqs. 10, 20, and 25 in Owocki et al. 2016) as a function of the polar coordinates (r, θ) of the upflowing, shocked, and downflowing plasma as a function of v_∞ , R_* , M_* , and $\dot{M}_{B=0}$.³ Since the P-Cygni absorption originates in the upflowing and downflowing plasma in the magnetosphere, we compare the densities derived from our spherical models with the sum $\rho_u + \rho_d$. Evaluating this for a typical viewing angle of $\theta = 60^\circ$ at different radii between 1.1 and $2.0 R_*$, and assuming $v_\infty = 1700 \text{ km s}^{-1}$ and the parameters derived in our model, we obtain values for $\log \dot{M}_{B=0}$ that range between -9.2 and $-8.8 [M_\odot \text{ yr}^{-1}]$. This is smaller than the values predicted by Vink et al. (2000) ($\approx -7.7 [M_\odot \text{ yr}^{-1}]$) by more than an order of magnitude, but in line with values reported for other stars of similar spectral type (e.g., Marcolino et al. 2009).

Since our radiative transfer assumes spherical symmetry, our result should be considered as an order-of-magnitude estimate (± 0.5 dex), and not as an accurate derivation of $\log \dot{M}_{B=0}$. From $\log \dot{M}_{B=0} \approx -9.0 [M_\odot \text{ yr}^{-1}]$ and $B_d \gtrsim 2 \text{ kG}$, and assuming again $v_\infty = 1700 \text{ km s}^{-1}$, we can constrain the Alfvén radius to $r_A \gtrsim 12 R_*$. From this, we can then estimate the true mass-loss from the star to $\log \dot{M} \lesssim -10.2 [M_\odot \text{ yr}^{-1}]$ (see equations 1-3 of Petit et al. 2017). However, we cannot exclude the possibility that some mass is lost from the star as a hot, X-ray emitting wind.

Importantly, while the magnetic field suppresses mass-loss from the star, it causes a density enhancement around it that leaves a spectroscopic signature. Therefore, the analysis of massive magnetic stars using appropriate models may enable one, in principle, to predict the mass-loss rates of non-magnetic stars of a similar

³ The model also introduces the parameter δ which describes the location of onset of matter infall along the dipole loop. The impact of this parameter is, for our estimation, negligible.

spectral type, which cannot be measured otherwise. Magnetic stars can therefore help to resolve the weak-wind problem.

3.4. The effect of X-rays on the stellar wind

The X-rays present in the stellar atmosphere are expected to affect the ionization balance via K-shell Auger ionization (Auger 1923). This effect is known to lead to high ionization stages such as N v and O vi (Cassinelli & Olson 1979; Oskinova et al. 2011). Moreover, X-rays affect diagnostic wind lines such as the resonance lines of C iv and Si iv. In our case, the N v $\lambda\lambda$ 1239, 1243 resonance lines are clearly present in the *HST* data, while our model predicts that N v is not significant in the stellar wind when not including X-rays. This suggests that X-rays contribute to the appearance of the UV spectrum.

Auger ionization via X-rays is accounted for in our model by assuming optically thin, thermally emitting filaments of shocked plasma embedded in the wind (Baum et al. 1992). While the topology of the X-ray emitting plasma may be more complex in reality, the important quantity here is merely the amount of ionizing radiation. The observed X-ray spectrum is characterized by two temperature components $T_{X,1}$ and $T_{X,2}$ and corresponding filling factors $X_{\text{fill},1}$ and $X_{\text{fill},2}$. These parameters are adopted from our analysis presented in Sect. 5, where only the two dominant components are accounted for. The onset radius of these filaments is set to $R_0 = 1.05 R_*$. Slightly different values deliver similar results, but onset radii which are too large trivially do not affect the spectra.

In Fig. 10, we compare the *HST* observations of the N v $\lambda\lambda$ 1239, 1243 and Si iv $\lambda\lambda$ 1394, 1403 resonance lines to our best-fitting model with and without X-rays. As discussed above, the N v resonance doublet is clearly seen in the observations but not in the model without X-rays. With the inclusion of X-rays, the N v lines appear. The N v line profiles have a similar shape to the C iv and Si iv resonance lines, suggesting that they too form in the magnetosphere. The strength of the lines implies a nitrogen abundance slightly lower than derived from our photospheric analysis by a factor of ≈ 0.7 , but this is within our given uncertainty. The X-rays also significantly influence the resonance lines of Si iv (see Fig. 10, right panel) and C iv. Including X-rays, our model successfully reproduces the main characteristics of our observations.

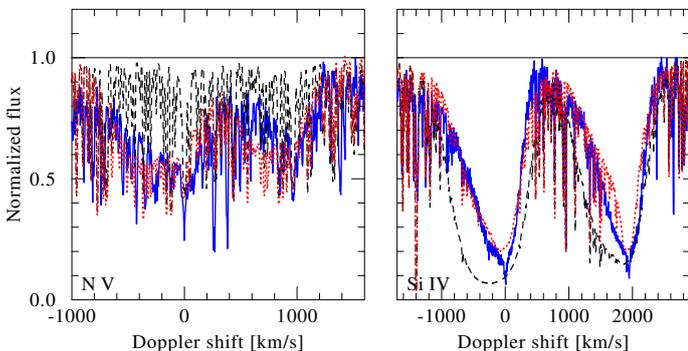


Fig. 10: Normalized *HST* observations of HD 54879 (blue solid line) compared to our best-fitting model (red dotted line) and to the same model without the inclusion of X-rays (black dashed line).

4. Evolutionary status

Using the parameters given in Table 1, we can derive an evolutionary scenario for HD 54879. Here, we use the BONNSAI⁴ tool (Schneider et al. 2014), which implements Bayesian statistics on a set of evolutionary tracks for massive stars (Brott et al. 2011) to constrain the best-fitting evolutionary channel. Previously performed by C2015, we repeat this procedure in light of the different parameters derived in this study. The parameters used are $\log L$, $\log T$, and $\log g_*$. The errors for $\log T_*$ and $\log g_*$, and $\log L$ are taken from Table 1. The peculiar C and N abundances are ignored.

Our BONNSAI solution predicts an initial mass of $M_{\text{ini}} = 16 \pm 1 M_{\odot}$ and an age of 5 ± 1 Myr. Our results are consistent with the findings by C2015 within errors. We do not find any indication that a merger event was involved in the formation of HD 54879, such as rejuvenation or rapid rotation. Nevertheless, traces for a past merger event may be unobservable if it occurred several million years ago.

We note that the evolution models used by the BONNSAI tool do not include the effect of magnetic fields. Thus, in addition to the uncertainties involved in evolution models of massive stars (e.g., overshooting, mass-loss rates), systematic errors due to the omission of magnetic fields in the models may interfere with our results. For example, according to our results in Sect. 3.3, HD 54879 loses significantly less mass than assumed by the BONNSAI evolution tracks. Because the magnetic field is known to strongly dampen the surface rotation, and in light of the very low $v \sin i$ value measured in this work, we used tracks which neglect rotationally induced mixing by setting the equatorial rotational velocity to zero ($v_{\text{rot, ini}} = 0 \text{ km s}^{-1}$).

5. X-ray spectral analysis

To analyze the *XMM-Newton* spectra, we used the standard spectral fitting software *XSPEC* (Arnaud 1996). The abundances were adopted from our spectral analysis in Sect. 3 (cf. Table 2). The X-ray flux of HD 54879 in the 0.2–12 keV band measured by the *XMM-Newton* is $\approx 2 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ (see Table 3). The unabsorbed flux corresponds to an X-ray luminosity at the distance of 1.0 kpc $\log L_X \approx 32.0 \text{ erg s}^{-1}$, resulting in $\log L_X/L_{\text{bol}} \approx -6$ (cf. Table 1). This is a higher value compared to other stars with similar spectral types. For example, the X-ray luminosity and the temperature of the hottest plasma found in μ Col (O9.5 V), ζ Oph (O9.2I V), 10 Lac (O9 V), σ Ori (O9.5 V), and other non-magnetic late O-dwarfs are at least an order of magnitude lower than in HD 54879 (Oskinova et al. 2006; Waldron & Cassinelli 2007; Huenemoerder et al. 2012).

In contrast, all known magnetic O-dwarfs – θ^1 Orionis C (O7 Vp), Tr 16-22 (O8.5 V), HD 57682 (O9.5 V), τ Sco – have relatively hard X-ray spectra and X-ray luminosities $\gtrsim 10^{32} \text{ erg s}^{-1}$, similar to HD 54879 (Schulz et al. 2000; Nazé et al. 2014). Thus, it appears safe to conclude that there is a clear dichotomy in the X-ray properties among non-magnetic and magnetic O-dwarfs, with the latter being significantly more X-ray luminous and displaying harder X-ray spectra than the former (the situation may be more complex in case of OB-supergiants and Of?p stars, see recent reviews by Nazé et al. 2014; ud-Doula & Nazé 2016). On this basis, we attribute the relatively high X-ray luminosity of HD 54879 to its magnetic nature.

It order to establish the temperature of the X-ray emitting plasma, we have analyzed the observed low-resolution *XMM-*

⁴ The BONNSAI web-service is available at www.astro.uni-bonn.de/stars/bonnasai

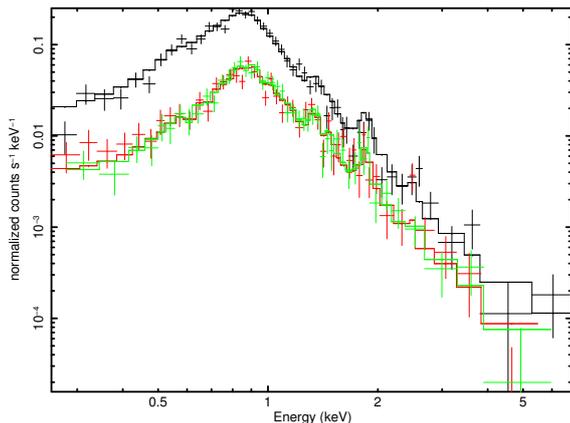


Fig. 11: *XMM-Newton* PN, MOS1, and MOS2 spectra of HD 54879 (black, red, and green curves, respectively) with error bars corresponding to 3σ with the best fit three-temperature thermal model (solid lines). The model parameters are shown in Table 1.

Newton spectra of HD 54879 (high-resolution RGS spectra have an insufficient signal and are not useful). The X-ray spectra of magnetic hot stars are usually well described by multi-temperature thermal plasma models (e.g., Oskinova et al. 2011; Nazé et al. 2014). This is also the case for HD 54879.

As a first step, we fitted the observed spectra in the 0.2–10.0 keV band with a thermal two-temperature spectral model that assumes optically thin plasma in collisional equilibrium. At this step, the absorption was modeled as originating in the cold interstellar medium (Wilms et al. 2000). The resulting spectral fit is statistically significant, with a reduced $\chi^2 = 1.3$. The best fit parameters are shown in Table 3.

There is no physical motivation for restricting the plasma temperature distribution to two temperatures. In fact, spectroscopic analyses show a continuous temperature distribution in hot stars (Wojdowski & Schulz 2005). Therefore, as a next step, we fitted to the observed EPIC spectra of HD 54879 a three-temperature spectral model. Also in this case, the fit is statistically significant, with $\chi^2 = 1.1$. The three-temperature model fit is shown in Fig. 11, and the associated parameters are shown in Table 3. The three temperature fit reveals the hottest plasma component with a temperature of about 2 keV (~ 20 MK). Better SNR is required to confirm the presence of this high-temperature plasma. We also tried to fit a four-temperature plasma model. However, the quality of the data is not sufficient to further constrain the temperature distribution.

In the framework of the magnetically confined wind shocks (MCWS) model (Babel & Montmerle 1997; ud-Doula & Owocki 2002), the wind plasma streams collide at the magnetic equator and give rise to a shock that heats the plasma. Hence, the maximum temperature follows from a Rankine-Hugoniot condition and cannot exceed a value determined by the maximum streaming velocity. Motivated by the recent discovery of a non-thermal component in the X-ray emission of the magnetic star HR 7355 (Leto et al. 2017), we also fitted the observed spectra with a composite two temperature thermal plasma and a non-thermal, power law component. The model provides a fit of similar quality to the observed spectra ($\chi^2 = 1.2$ for 138 degrees of freedom). The best fit thermal plasma temperatures are $kT_1 = 0.18 \pm 1.2$ [keV], $kT_2 = 0.74 \pm 0.02$ [keV], and the power-law exponent is $\Gamma = 2.5 \pm 0.7$. Thus, the quality of the data is not

Table 3: The X-ray spectral parameters derived from the *XMM-Newton* observations of HD 54879 assuming a three-temperature plasma model.

Two temperature thermal model	
N_H [10^{21} cm $^{-2}$]	4.2 ± 0.4
kT_1 [keV]	0.18 ± 0.01
EM_1 [10^{54} cm $^{-3}$]	6.60 ± 0.4
kT_2 [keV]	0.73 ± 0.02
EM_2 [10^{54} cm $^{-3}$]	2.46 ± 0.03
$\langle kT \rangle \equiv \sum_i kT_i \cdot EM_i / \sum_i EM_i$ [keV]	0.33
Three temperature thermal model	
N_H [10^{21} cm $^{-2}$]	3.6 ± 0.5
kT_1 [keV]	0.18 ± 0.01
EM_1 [10^{54} cm $^{-3}$]	3.70 ± 0.01
kT_2 [keV]	0.74 ± 0.02
EM_2 [10^{54} cm $^{-3}$]	2.01 ± 0.05
kT_3 [keV]	2.1 ± 0.6
EM_3 [10^{54} cm $^{-3}$]	0.31 ± 0.01
$\langle kT \rangle \equiv \sum_i kT_i \cdot EM_i / \sum_i EM_i$ [keV]	0.46
Flux [10^{-13} erg cm $^{-2}$ s $^{-1}$]	1.6
L_X^a [erg s $^{-1}$]	1×10^{32}
$\log L_X / L_{\text{bol}}$	-6

Notes. ^(a) corrected for the interstellar absorption in the 0.2–12 keV band

sufficient to discriminate between three-temperature thermal and two-temperature thermal plus non-thermal component models.

Interestingly, the total absorption derived from our X-ray analysis is found to be somewhat higher than that derived from the analysis of UV and optical data (Sect. 3). Using the averaged relation $N_H = E_{B-V} \cdot 6 \cdot 10^{21}$ cm $^{-2}$ (Gudennavar et al. 2012), $E_{B-V} = 0.35$ implies $N_{H,\text{ISM}} = 2.1 \times 10^{21}$ cm $^{-2}$, which is lower by almost a factor of two compared to what is found from the X-ray data (see Table 3). Since the absorption derived from X-rays includes both interstellar and intrinsic absorption, it is possible that the additional X-ray absorption occurs in the magnetosphere of HD 54879. Alternatively, this discrepancy could be merely a consequence of the large error on N_H (see Table 3) when derived from our low resolution X-ray spectra, and of the uncertain relations between N_H , E_{B-V} , and the extinction parameters A_V and R_V (e.g., Güver & Özel 2009).

To test whether some X-ray absorption originates in the magnetosphere, we searched for the presence of absorption edges corresponding to the leading ionization stages of abundant metals in HD 54879. Fitting the observed spectra with an absorption model that includes multiple ions did not statistically improve the fit. Hence, we concentrated on searching absorption edges of individual ions. Fitting a three-temperature plasma model that, in addition to the interstellar absorption, is also attenuated by a “warm” absorber resulted in a detection of an absorption edge at 0.61 ± 0.15 keV with an optical depth 0.9 ± 0.6 . This energy corresponds to the absorption edges of O III-VII (Verner & Yakovlev 1995). The fact that these are leading ionization stages of oxygen in the magnetosphere of HD 54879 provides additional support for the detection of the absorption edge. Including ionization edges in the X-ray spectral model does not

significantly affects the derived value of the total ISM absorption $N_{\text{H,ISM}}$. Hence, while the data agree with some X-ray absorption due to the warm material trapped in the magnetosphere, the quality of our data prevents us from measuring the warm absorption component unambiguously.

In Sect. 3.3, we showed that the true wind parameters (\dot{M} , v_∞) cannot be derived for this object based on UV/optical spectra. However, as suggested by Lucy (2012), it is possible that a significant amount of matter leaves the star in a very hot phase and emits in X-rays. From high resolution X-ray spectra, one could study in detail the properties of the stellar wind of the HD 54879, if it indeed exists. To shed light on the urgent weak wind problem, as well as on the true mass-loss rates of massive main sequence stars, we therefore encourage future observational campaigns to acquire high resolution X-ray spectra for HD 54879.

6. Variability

6.1. Photometric variability

In Fig. 12, we present a lightcurve collected by the All Sky Automated Survey 3 (ASAS3, six-pixel wide aperture) for our target (Pojmanski 2002). The dataset was scanned for significant periods using Fourier and F-statistics, but none could be identified. Striking are the sudden decreases in brightness, which only last for ≈ 1 d, evident in Fig. 12. We checked whether these events occur periodically and can safely reject this possibility, excluding binary eclipses as a possible explanation. Sudden brightness changes could also arise from outbursts. However, the intensity of change in magnitude and the short time scale over which these events occur do not plausibly agree with such a model. Similar outliers can be seen in other lightcurves obtained by the ASAS (e.g., Paczyński et al. 2006). We conclude that these events are most likely observational artefacts.

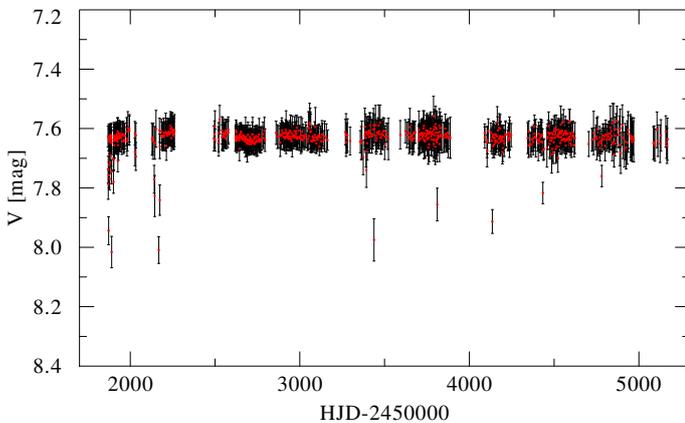


Fig. 12: ASAS3 lightcurve of HD 54879

6.2. Spectroscopic variability

HD 54879 appears to exhibit both small-scale and large-scale variability that are especially evident in the $\text{H}\alpha$ line (see Fig. 3). While the dynamic structure of the magnetosphere implies some stochastic variability (e.g., Sundqvist et al. 2012; ud-Doula et al. 2013), it is also expected that a periodic variability that correlates with the rotational period of the star should be present (e.g., Stibbs 1950; Townsend & Owocki 2005; Wade et al. 2011).

To illustrate the short scale variability, we refer the reader to Fig. 3 in Sect. 3. The figure shows that, while $\text{H}\alpha$ maintains a

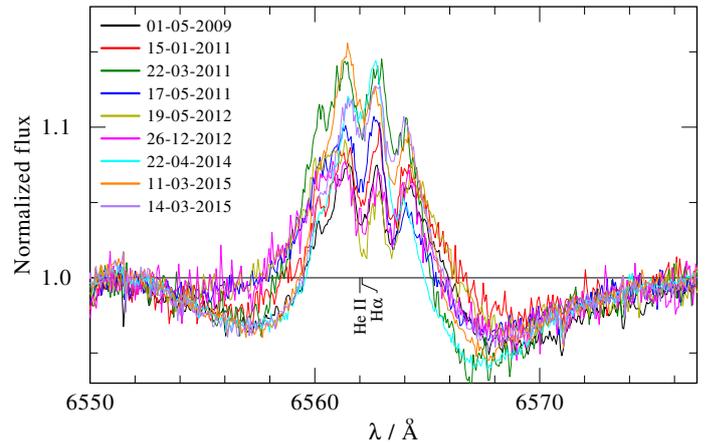


Fig. 13: FEROS, FIES, and HARPS spectra of $\text{H}\alpha$ taken in the years 2009 – 2015 (see legend and Table A.1)

similar equivalent width between 11-03-2015 and 14-03-2015, some variability can be seen in the strength of the leftmost emission peak and the right absorption wing. In other words, the star exhibits some variability on the scale of days, which can be considered short-scale in our context. The third spectrum, taken about half a year prior to the others (April 22, 2014), shows a higher equivalent width (less emission) and a narrower profile. This suggests that a mechanism is present that is responsible for a long-term variability. Fig. 13 shows nine spectra of $\text{H}\alpha$ taken in the years 2009 – 2015 (see Table A.1 and legends). One can see clear changes in the equivalent widths and full width half maxima (FWHM) of the profiles, as well as differing strengths of the absorption wings.

These spectra alone are not sufficient to determine a period. For this reason, we acquired 35 amateur spectra (P. Luckas, priv. comm.) collected specifically for this project (see Sect. 2). The spectra are of lower quality, but offer the advantage of a much denser time coverage. A few selected spectra are shown in Fig. 14.

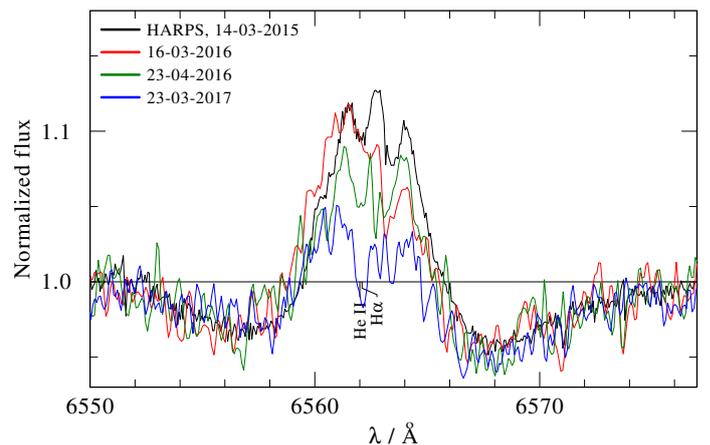


Fig. 14: A few amateur SPO spectra of $\text{H}\alpha$, shown with a single HARPS spectrum for comparison

It is not trivial to identify the quantity whose time series would best constrain a period for the global variability. A rotational modulation can lead to partial polar/equatorial view (relative to the magnetic axis) of the disk-like, $\text{H}\alpha$ emitting structure, which in turn leads to a periodic variation in the equivalent widths (e.g., Sundqvist et al. 2012; ud-Doula et al. 2013). We therefore

measured the equivalent width of the $H\alpha$ line in the range 6551 – 6576 Å in all available spectra. The error bars are attributed mostly to uncertainty in rectification, and are roughly estimated by $\Delta\lambda / \text{SNR}$, where $\Delta\lambda = 25$ Å is the integration domain. The results are listed in Table A.1 and shown in Fig. 15.

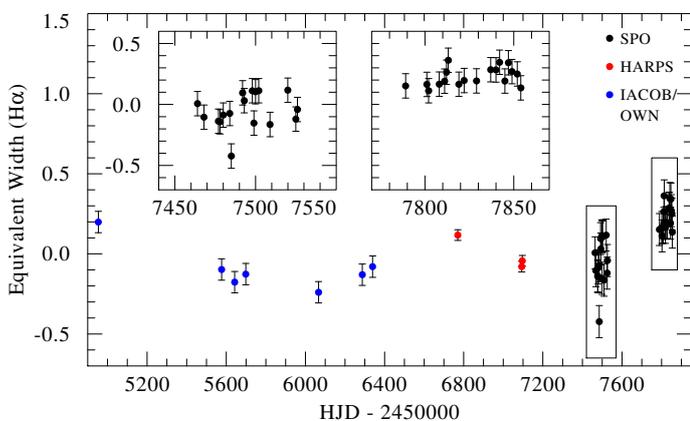


Fig. 15: Equivalent widths of $H\alpha$ versus Heliocentric Julian date, as measured in the IACOB/OWN, HARPS and amateur SPO spectra

Figure 15 seems to be suggestive of a modulation of the equivalent widths over a very long period, of the order of years. A Fourier analysis of the signal (Scargle 1982; Horne & Baliunas 1986) suggests a period of ≈ 5 yr as the only significant period. This result is consistent with the upper limit of $v \sin i < 4$ km s $^{-1}$ derived in Sect. 3. From the derived period and stellar radius, and assuming the probable value $i \approx 60^\circ$, this suggests a very low rotational velocity of $v_{\text{eq}} \approx 0.2$ km s $^{-1}$. Additional attempts to constrain the period using other proxies such as the line width or its centroid did not result in any coherent periods. We therefore suggest in this paper a rotational period of ≈ 5 yr, but more spectra will be necessary to validate our claim. Overall, this result is consistent with the fact that magnetic stars are very slow rotators that lost their angular momentum via magnetic braking (Weber & Davis 1967; ud-Doula & Owocki 2002).

7. Summary

We performed a comprehensive, multiwavelength analysis of HD 54879 (O9.7 V). With a dipole magnetic field of $B_d \gtrsim 2$ kG, it is one of about ten magnetized O-stars known. Using high quality X-ray, UV, and optical spectra acquired by the *XMM-Newton*, *HST*, and HARPS instruments, respectively, we could derive the X-ray properties of HD 54879 and analyze its atmosphere and wind. Moreover, 45 spectra were used to constrain a periodic variability of the $H\alpha$ line. We conclude the following:

- The fundamental stellar parameters (T_* , L , $\log g$) of HD 54879 are typical for its spectral type. Nitrogen and carbon abundances are found to be sub-solar.
- The projected rotational velocity $v \sin i$, the microturbulent velocity ξ_{ph} , and the macroturbulent velocity v_{mac} , are all found to be smaller than 4 km s $^{-1}$.
- The X-ray spectrum can be well-fitted with a thermal model accounting for either two or three components, and implies a higher-than-average X-ray luminosity ($\log L_X/L_{\text{bol}} = -6.0$). In the three-component model, the X-ray temperature reaches values up to $T_X \approx 20$ MK.

- Variability of $H\alpha$ equivalent widths is suggestive of a very long period of the order of ≈ 5 yr, consistent with the low $v \sin i$ value.
- The mass-loss rate that our target would have in the absence of a magnetic field could be roughly estimated to be $\log \dot{M}_{B=0} \approx -9.0$ [$M_\odot \text{ yr}^{-1}$]. This is significantly less than theoretically predicted (Vink et al. 2000), but is in line with what is found for non-magnetic stars of similar spectral types (Marcolino et al. 2009). With an Alfvén radius of $r_A \gtrsim 12 R_*$, the true mass-loss rate of HD 54879 is estimated as $\log \dot{M} \lesssim -10.2$ [$M_\odot \text{ yr}^{-1}$].

To conclude, we would like to point out that, as our study illustrates, slowly rotating magnetic stars can provide important constraints on the *weak wind problem*. The spectra of main sequence OB-type stars often exhibit very little or no signatures for a stellar wind in UV spectra, which contain the principle diagnostics for mass-loss rates. Magnetic fields that confine the stellar winds enhance the circumstellar densities and result in clear spectral features. Using sophisticated models for the magnetospheres of massive magnetic stars, one can therefore infer values for $\log \dot{M}_{B=0}$. These should be similar to *true* mass-loss rates for *non-magnetic* main sequence stars of the same spectral type.

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Appendix A: Observations and H α equivalent widthsTable A.1: Compilation of optical observations and measured H α equivalent widths for HD 54879.

No.	Instrument	Date	HJD	EW _{Hα} [Å]
1	FEROS	01-05-2009	2454953.54	-0.11
2	FIES	15-01-2011	2455576.54	-0.17
3	FEROS	22-03-2011	2455642.54	-0.45
4	FEROS	17-05-2011	2455698.54	-0.27
5	FEROS	19-05-2012	2456066.54	-0.46
6	FIES	26-12-2012	2456287.54	-0.17
7	FIES	16-02-2013	2456339.54	-0.12
8	HARPS	22-04-2014	2456770.04	0.02
9	HARPS	11-03-2015	2457093.04	-0.09
10	HARPS	14-03-2015	2457096.04	-0.18
11	SPO	16-03-2016	2457464.02	0.00
12	SPO	20-03-2016	2457468.02	-0.10
13	SPO	29-03-2016	2457476.99	-0.13
14	SPO	30-03-2016	2457477.98	-0.14
15	SPO	01-04-2016	2457479.97	-0.08
16	SPO	05-04-2016	2457483.98	-0.07
17	SPO	06-04-2016	2457484.97	-0.42
18	SPO	13-04-2016	2457491.96	0.09
19	SPO	14-04-2016	2457493.01	0.03
20	SPO	19-04-2016	2457497.96	0.11
21	SPO	20-04-2016	2457498.97	-0.15
22	SPO	21-04-2016	2457499.97	0.10
23	SPO	23-04-2016	2457501.97	0.11
24	SPO	30-04-2016	2457508.94	-0.16
25	SPO	11-05-2016	2457519.96	0.11
26	SPO	16-05-2016	2457524.93	-0.11
27	SPO	17-05-2016	2457525.94	-0.04
28	SPO	04-02-2017	2457789.07	0.15
29	SPO	16-02-2017	2457801.03	0.16
30	SPO	17-02-2017	2457802.03	0.11
31	SPO	23-02-2017	2457808.02	0.16
32	SPO	26-02-2017	2457811.03	0.19
33	SPO	27-02-2017	2457812.09	0.26
34	SPO	28-02-2017	2457813.11	0.36
35	SPO	06-03-2017	2457819.07	0.16
36	SPO	09-03-2017	2457822.02	0.19
37	SPO	16-03-2017	2457829.02	0.19
38	SPO	24-03-2017	2457837.00	0.28
39	SPO	27-03-2017	2457839.99	0.28
40	SPO	29-03-2017	2457841.99	0.34
41	SPO	01-04-2017	2457844.99	0.19
42	SPO	03-04-2017	2457846.98	0.34
43	SPO	05-04-2017	2457848.98	0.26
44	SPO	08-04-2017	2457851.97	0.24
45	SPO	10-04-2017	2457853.97	0.13